

$$8.2 \quad \psi(x) = e^{if(x)/\hbar}$$

$$\frac{d\psi}{dx} = \frac{i}{\hbar} f'(x) e^{if(x)/\hbar}$$

$$\frac{d^2\psi}{dx^2} = \left(-\frac{[f'(x)]^2}{\hbar^2} + \frac{i}{\hbar} f''(x) \right) e^{if(x)/\hbar}$$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$

$$\frac{d^2\psi}{dx^2} = -\frac{p^2}{\hbar^2} \psi$$

$$-\frac{f_1'^2}{\hbar^2} + \frac{i}{\hbar} f'' = -\frac{p^2}{\hbar^2}$$

$$i\hbar f'' - (f')^2 + p^2 = 0$$

$$(b.) \quad f(x) = f_0(x) + \hbar f_1(x) + \hbar^2 f_2(x) + \dots + \hbar^n f_n(x)$$

$$f'(x) = f_0'(x) + \hbar f_1'(x) + \hbar^2 f_2'(x) + \dots$$

$$f''(x) = f_0''(x) + \hbar f_1''(x) + \hbar^2 f_2''(x) + \dots$$

$$(f_1')^2 = (f_0'(x))^2 + 2\hbar f_0' f_1' + \hbar^2 (2f_0' f_2' + f_1'^2) + 2\hbar^3 [f_0' f_3' + f_1' f_2'] + O(\hbar^4)$$

$$i\hbar f'' = i(\hbar f_0'' + \hbar^2 f_1'' + \hbar^3 f_2'' + O(\hbar^4))$$

$$(i\hbar f_0'' + i\hbar^2 f_1'' + i\hbar^3 f_2'') - [f_0'^2 + 2\hbar f_0' f_1' + \hbar^2 (2f_0' f_2' + f_1'^2) + 2\hbar^3 (f_0' f_3' + f_1' f_2')] + p^2 = 0$$

Grouping like powers of \hbar :

$$\hbar^0: \quad -f_0'^2 + p^2 = 0 \rightarrow (f_0')^2 = p^2$$

$$\hbar^1: \quad i f_0'' - 2f_0' f_1' = 0 \rightarrow i f_0'' = 2f_0' f_1'$$

$$\hbar^2: \quad i f_1'' - 2f_0' f_2' - (f_1')^2 = 0 \rightarrow i f_1'' = 2f_0' f_2' + (f_1')^2$$

$$\hbar^3: \quad i f_2'' - 2(f_0' f_3' + f_1' f_2') = 0 \rightarrow i f_2'' = 2f_0' f_3' + 2f_1' f_2'$$

$$(c.) \quad (f_0')^2 = p^2 \rightarrow f_0' = \pm p$$

$$f_0 = \pm \int p(x) dx + C \quad (2)$$

$$i f_0'' = 2 f_0' f_1'$$

$$\frac{i}{2} \frac{f_0''}{f_0'} = f_1' \rightarrow f_1' = \frac{i}{2} \left(\frac{\pm p'}{\pm p} \right) = \frac{i}{2} \left(\frac{\frac{d}{dx} p}{p} \right)$$

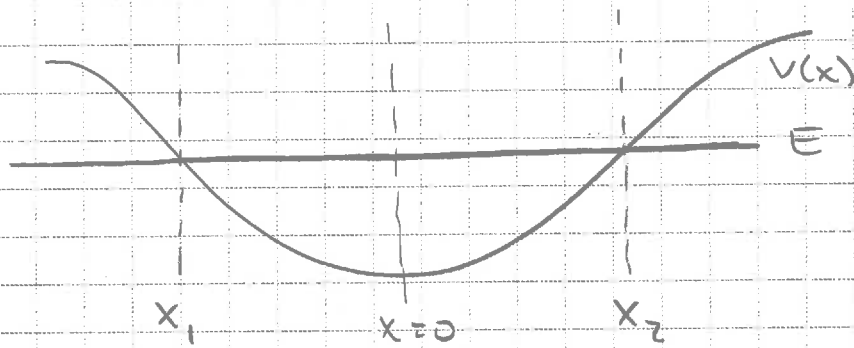
$$f_1' = \frac{i}{2} \frac{d}{dx} \left(\frac{dp}{p} \right) \Rightarrow \frac{i}{2} \frac{d}{dx} [\ln(p)] \Rightarrow f_1 = \frac{i}{2} \ln p + C$$

$$\psi = \exp\left(\frac{i f(x)}{\hbar}\right) = \exp\left[\underbrace{\frac{i}{\hbar} \int p(x) dx}_{f_0} + \underbrace{\frac{i \hbar}{2} \ln(p) + C_0}_{f_1}\right]$$

$$\psi = \exp\left(\frac{\pm i}{\hbar} \int p(x) dx\right) \exp\left(-\frac{1}{2} \ln(p)\right) C$$

$$\psi = \frac{C}{\sqrt{p}} \exp\left[\frac{\pm i}{\hbar} \int p(x) dx\right] \quad \text{QED}$$

8.7



We need to use the potential well - /
no wells:

$$\int_{x_1}^{x_2} p(x) dx = \pi \hbar \left(n - \frac{1}{2}\right) \quad \text{Eqn. 8.5}$$

$$p(x) = \sqrt{2m \left(E - \frac{1}{2} m \omega^2 x^2\right)}$$

$$E = \frac{1}{2} m \omega^2 x_1^2, \quad \text{where } x_1 = -x_2$$

$$-\sqrt{\frac{2E}{m\omega^2}} = x_1$$

$$\sqrt{\frac{2E}{m\omega^2}} = x_2$$

$$(n-\frac{1}{2})\pi\hbar = \int_{x_1}^{x_2} p(x) dx = \int_{x_1}^{x_2} \sqrt{2m \left[\frac{2E}{m\omega^2} - x^2 \right]^{1/2}} dx \quad (3)$$

$$= 2m\omega \int_0^{x_2} \sqrt{x_2^2 - x^2} dx, \text{ where } x_2 \equiv \sqrt{\frac{2E}{m\omega^2}}$$

$$= m\omega \left(x\sqrt{x_2^2 - x^2} + x_2^2 \sin^{-1}\left(\frac{x}{x_2}\right) \right) \Big|_0^{x_2} = m\omega \left[0 + x_2^2 \sin^{-1}(1) + \cancel{x_2^2 \sin^{-1}(0)} \right]$$

$$\sin^{-1}(1) = \frac{\pi}{2} \quad \sin^{-1}(0) = 0$$

$$(n-\frac{1}{2})\pi\hbar = m\omega x_2^2 \frac{\pi}{2}$$

$$= \frac{1}{2} m\omega \frac{2E}{m\omega^2} = \frac{E}{\omega}$$

$$\hbar\omega(n-\frac{1}{2}) = E \quad \text{for } n=1, 2, 3, \dots$$

But since $n-\frac{1}{2}$ for $n=1$ is $\frac{1}{2}$, we can switch to $E = \hbar\omega(n+\frac{1}{2})$ for $n=0, 1, 2, \dots$ instead, which is exactly the result for the HO.

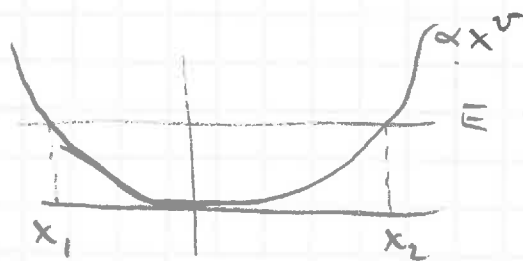
8.11 Again, we need to use Eqn. 8.57

$$(n-\frac{1}{2})\pi\hbar = \int_{x_1}^{x_2} p(x) dx$$

Since we're dealing w/ an abs. value in our potential, we know that $p(x)$ is even so:

$$(n-\frac{1}{2})\pi\hbar = 2 \int_0^{x_2} p(x) dx$$

$$p(x) = \sqrt{2m(E - \alpha x^{\nu})}$$



$$E = \alpha x_2^{\nu}$$

$$z \equiv \frac{\alpha}{E} x^{\nu}$$

$$x = \left(\frac{2E}{\alpha}\right)^{1/2} \quad dx = \left(\frac{E}{\alpha}\right)^{1/2} \frac{1}{\sqrt{v}} z^{v-1} dz \quad (4)$$

$$(n - \frac{1}{2})\pi\hbar = 2\sqrt{2mE} \int_0^{x_2} \left(1 - \frac{\alpha}{E} x^2\right)^{1/2} dx$$

$$= 2\sqrt{2mE} \left(\frac{E}{\alpha}\right)^{1/2} \frac{1}{\sqrt{v}} \int_0^1 z^{1/2-1} \sqrt{1-z} dz = 2\sqrt{2mE} \left(\frac{E}{\alpha}\right)^{1/2} \frac{\Gamma(1/2)\Gamma(3/2)}{\Gamma(1/2+3/2)}$$

$$= 2\sqrt{2mE} \left(\frac{E}{\alpha}\right)^{1/2} \frac{\Gamma(1/2+1) \frac{1}{2}\sqrt{\pi}}{\Gamma(1/2+3/2)} = \sqrt{2\pi mE} \left(\frac{E}{\alpha}\right)^{1/2} \frac{\Gamma(1/2+1)}{\Gamma(1/2+3/2)}$$

$$(n - \frac{1}{2})\pi\hbar = \sqrt{2\pi m} \left(\frac{1}{\alpha}\right)^{1/2} \frac{\Gamma(1/2+1)}{\Gamma(1/2+3/2)} E^{1/2+1/2}$$

$$\frac{(n - \frac{1}{2})\pi\hbar}{\sqrt{2\pi m}} \alpha^{1/2} \frac{\Gamma(1/2+3/2)}{\Gamma(1/2+1)} = E^{1/2+1/2}$$

$$\therefore E_n = \alpha \left[(n - \frac{1}{2})\hbar \sqrt{\frac{\pi}{2m\alpha}} \frac{\Gamma(1/2+3/2)}{\Gamma(1/2+1)} \right]^{\frac{2v}{v+2}}$$

$$v = 2$$

$$E_n = \alpha (n - \frac{1}{2})\hbar \sqrt{\frac{\pi}{2m\alpha}} \frac{\Gamma(2)}{\Gamma(3/2)}$$

$$= \alpha (n - \frac{1}{2})\hbar \frac{2}{\sqrt{\pi}} \sqrt{\frac{\pi}{2m\alpha}}$$

$$E_n = (n - \frac{1}{2})\hbar \sqrt{\frac{2\alpha}{m}}$$

For the HO, $\alpha = \frac{1}{2}m\omega^2$, so: $E_n = (n - \frac{1}{2})\hbar \sqrt{\frac{2(\frac{1}{2}m\omega^2)}{m}}$

$$E_n = (n - \frac{1}{2})\hbar\omega \quad \checkmark$$

$$1 \Rightarrow \Gamma(n) = (n-1)!$$

$$\Gamma(\alpha) = (\alpha-1)\Gamma(\alpha-1)$$

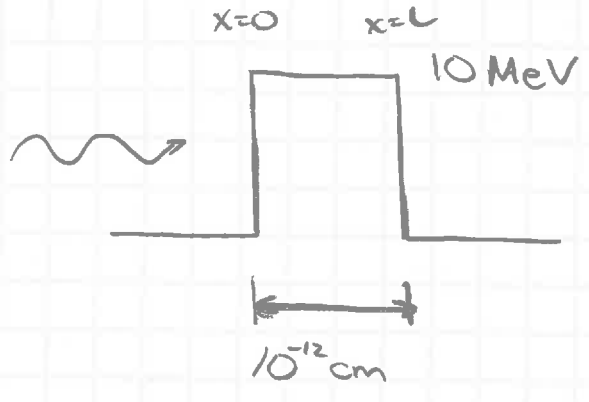
$$\Gamma(3/2) = \frac{1}{2}\Gamma(1/2)$$

$$\Gamma(1/2) = \sqrt{\pi} \quad \text{so:}$$

$$\Gamma(3/2) = \frac{\sqrt{\pi}}{2}$$

4.

$KE = 4 \text{ MeV}$
 $L = 10^{-12} \text{ cm}$
 $V_0 = 10 \text{ MeV}$



$$\Psi_I(x) = A_0 e^{ikx} + A e^{-ikx}$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\Psi_{II}(x) = B e^{-\alpha x} + C e^{+\alpha x}$$

$$\alpha = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

$x=0$

① $A_0 + A = B + C$ (continuity of $\Psi(x)$)

② $ik(A_0 - A) = -\alpha(B - C)$ (continuity of $\Psi'(x)$)

$x=L$

③ $D e^{ikL} = B e^{-\alpha L} + C e^{+\alpha L}$ (continuity of $\Psi(x)$)

④ $ik D e^{ikL} = -\alpha(B e^{-\alpha L} - C e^{+\alpha L})$ (continuity of $\Psi'(x)$)

① $A_0 + A = B + C$

② $\oplus A_0 + A = \frac{-\alpha}{ik}(B - C)$

$$2A_0 = B\left(1 - \frac{\alpha}{ik}\right) + C\left(1 + \frac{\alpha}{ik}\right)$$

$$2ikA_0 = B(ik - \alpha) + C(ik + \alpha)$$

③ $D e^{ikL} - C e^{+\alpha L} = B e^{-\alpha L}$

④ $\frac{ik}{-\alpha} D e^{ikL} = B e^{-\alpha L} - C e^{+\alpha L}$

Putting ③ into ④

$$-\frac{ik}{\alpha} D e^{ikL} = D e^{ikL} - C e^{+\alpha L} - C e^{+\alpha L}$$

$$-(1 + \frac{ik}{\alpha})De^{ikL} = -2Ce^{+\alpha L}$$

$$\frac{1}{2\alpha}(\alpha + ik)De^{ikL} = Ce^{+\alpha L}$$

$$\frac{1}{2\alpha}(\alpha + ik)De^{(ik-\alpha)L} = C$$

$$\textcircled{3} \quad De^{ikL} - Ce^{+\alpha L} = Be^{-\alpha L}$$

$$De^{ikL}(1 - \frac{1}{2\alpha}(\alpha + ik)) = Be^{-\alpha L}$$

$$De^{(k+\alpha)L} [1 - \frac{1}{2\alpha}(\alpha + ik)] = B$$

$$2ikA_0 = B(ik - \alpha) + C(ik + \alpha)$$

$$2ikA_0 = De^{ikL} \left\{ \left[e^{\alpha L} \left[1 - \frac{1}{2\alpha}(\alpha + ik) \right] \right] (ik - \alpha) + \left[\frac{1}{2\alpha}(\alpha + ik)^2 \right] e^{-\alpha L} \right\}$$

$$4i\alpha k A_0 = De^{ikL} \left\{ (\alpha - ik)^2 e^{\alpha L} + (\alpha + ik)^2 e^{-\alpha L} \right\}$$

$$\frac{D}{A_0} = \frac{4i\alpha k e^{-ikL}}{-(\alpha - ik)^2 e^{\alpha L} + (\alpha + ik)^2 e^{-\alpha L}}$$

Let's compute αL first: $\alpha L = \frac{\sqrt{2m(V_0 - E)}}{\hbar} L$

$$\alpha_p L = \frac{\sqrt{2 \times 1.67 \times 10^{-27} \text{ kg} (6 \text{ MeV}) \times 1.602 \times 10^{-13} \text{ J}}}{1.055 \times 10^{-34} \text{ J}\cdot\text{s}} \quad m_p = 3.34 \times 10^{-27} \text{ kg}$$

$$m_p = 1.67 \times 10^{-27} \text{ kg}$$

$$\alpha_p L = 5.37 \quad (\text{bigger for deuteron})$$

IF $\alpha L \gg 1$, then:

$$\frac{D}{A_0} \approx \frac{4i\alpha k e^{-ikL}}{-(\alpha - ik)^2 e^{\alpha L}}$$

$$\text{so } T = \left| \frac{D}{A_0} \right|^2 \approx \frac{16\alpha^2 k^2}{(\alpha^2 + k^2)^2} e^{-2\alpha L}$$

$$T_p = \frac{16\alpha_p^2 k_p^2}{(\alpha_p^2 + k_p^2)^2} e^{-2(5.37)}$$

$$\alpha_p = \frac{\sqrt{2m_p(V_0 - E)}}{\hbar} = 5.37 \times 10^{+14} \text{ cm}^{-1}$$

$$k_p = \frac{\sqrt{2m_p E}}{\hbar} = 4.39 \times 10^{+14} \text{ cm}^{-1}$$

$$T_p \approx \frac{16 (5.37 \times 10^{14} \text{ cm}^{-1})^2 (4.39 \times 10^{14} \text{ cm}^{-1})^2}{[(5.37 \times 10^{14} \text{ cm}^{-1})^2 + (4.39 \times 10^{14} \text{ cm}^{-1})^2]^2} e^{-10.74}$$

3.84

$$T_p \approx 8.32 \times 10^{-5}$$

$$T_D \approx \frac{16 \alpha_D^2 k_D^2}{(\alpha_D^2 + k_D^2)^2} e^{-2\alpha_D}$$

$$\alpha_D = \frac{\sqrt{2m_0(V_0 - E)}}{\hbar} = 7.61 \times 10^{14} \text{ cm}^{-1}$$

$$k_D = \frac{\sqrt{2m_0 E}}{\hbar} = 6.22 \times 10^{14} \text{ cm}^{-1}$$

$$T_D = \frac{16 (7.61 \times 10^{14} \text{ cm}^{-1})^2 (6.22 \times 10^{14} \text{ cm}^{-1})^2}{[(7.61 \times 10^{14} \text{ cm}^{-1})^2 + (6.22 \times 10^{14} \text{ cm}^{-1})^2]^2} e^{-2(7.61)} = 9.4 \times 10^{-7} \approx 10^{-6}$$

3.84

5. (a) $V(x) = mgh = mgx$ for $x > 0$

$$\frac{-\hbar^2}{2m} \frac{d^2\psi}{dx^2} + mgx\psi = E\psi$$

$$\frac{d^2\psi}{dx^2} = -2m \left(\frac{E - mgx}{\hbar^2} \right) \psi$$

$$\frac{d^2\psi}{dx^2} = \frac{2m^2g}{\hbar^2} \left(x - \frac{E}{mg} \right) \psi$$

α^3 z

$$\frac{d^2\psi}{dz^2} = \alpha^3 z \psi$$

We can simplify this further (and put it into the Airy eqn form) by choosing $z = \alpha y$

$$\alpha^2 \frac{d^2\psi}{dz^2} = \alpha^3 z \psi$$

$$\frac{d^2\psi}{dz^2} = z \psi \quad (\text{Airy eqn.})$$

The general soln is: $\psi = a Ai(z) + b Bi(z)$

Since $\text{Bi}(z)$ blows up for large z , $b=0$.

$$\therefore \psi(x) = a \text{Ai}(z) = a \text{Ai}\left[\alpha\left(x - \frac{E}{mg}\right)\right], \text{ where}$$
$$\alpha = \left(\frac{2m^2g}{\hbar^2}\right)^{1/3}$$

(b.) $\psi(0) = 0$ since $V(0) = \infty$

\therefore At the zeroes of $\text{Ai}(z)$ we can find the allowed energies. $\text{Ai}(z) = 0$ @ $a_1 = -2.34$, $a_2 = -4.09$, and $a_3 = -5.52$.

$$\textcircled{1} \quad \alpha\left(-\frac{E_1}{mg}\right) = -2.34 \rightarrow \left(\frac{2m^2g}{\hbar^2}\right)^{1/3} \frac{E_1}{mg} = 2.34$$
$$E_1 = \left(\frac{\hbar^2 mg^2}{2}\right)^{1/3} 2.34$$

$$m = 0.1 \text{ kg}$$
$$g = 9.8 \text{ m/s}^2$$

$$E_1 = 8.81 \times 10^{-23} \text{ J}$$

$$\textcircled{2} \quad E_2 = \left(\frac{\hbar^2 mg^2}{2}\right)^{1/3} 4.09$$

$$E_2 = 1.54 \times 10^{-22} \text{ J}$$

$$\textcircled{3} \quad E_3 = \left(\frac{\hbar^2 mg^2}{2}\right)^{1/3} 5.52$$

$$E_3 = 2.08 \times 10^{-22} \text{ J}$$

$$\text{(c.) } E_1 = \left(\frac{\hbar^2 mg^2}{2}\right)^{1/3} 2.34 = 1.15 \times 10^{-13} \text{ eV}$$

$$2\langle T \rangle = \left\langle x \frac{dV}{dx} \right\rangle$$

$$\frac{dV}{dx} = \frac{d}{dx}(mgx) = mg$$

$$\left\langle x \frac{dV}{dx} \right\rangle = \langle mgx \rangle = \langle V \rangle$$

$$2\langle T \rangle = \langle x \frac{dV}{dx} \rangle = \langle V \rangle \rightarrow \langle T \rangle = \frac{1}{2} \langle V \rangle$$

(9)

$$\langle H \rangle = \langle T \rangle + \langle V \rangle = E_n$$

$$= \frac{1}{2} \langle v \rangle + \langle v \rangle = E_n$$

$$= \frac{3}{2} \langle v \rangle = E_n$$

$$= \frac{3}{2} mg \langle x \rangle = E_n$$

$$\langle x \rangle = \frac{2E_n}{3mg}$$

$$\langle x \rangle = \frac{2E_1}{3mg} = \frac{2(1.15 \times 10^{-13} \text{ eV} \times 1.602 \times 10^{-19} \text{ J/eV})}{3(9.11 \times 10^{-31} \text{ kg})(9.8 \text{ m/s}^2)} = 1.3747 \times 10^{-3} \text{ m}$$

$$\boxed{\langle x \rangle = 1.37 \text{ mm}}$$

